

# Resonance Frequency and Bandwidth of Rectangular Microstrip Antenna on Thick Substrate

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**Abstract**—We report an improved cavity model called modified Wolff model (MWM) to compute the resonance frequency of a rectangular patch antenna on the thick PTFE substrate ( $0.037\lambda_g$ – $0.229\lambda_g$ ). The MWM accounts for the effect of anisotropy and total losses on the resonance frequency, therefore it has a maximum deviation 2% against the experimental results. The previous cavity model, the multiport cavity model, and MOM based commercial software, Ensemble, compute resonance frequency with deviation between 4%–36%. Results on the bandwidth computed by these models have also been compared against the experimental results.

**Index Terms**—Microstrip antennas, rectangular antennas, thick patch.

## I. INTRODUCTION

THE RECTANGULAR microstrip antenna shown in Fig. 1 is normally designed on the electrically thin substrates of the order of  $h < 0.02\lambda_g$ , where the  $\lambda_g$  is the guided wavelength in the substrate of relative permittivity given by  $\lambda_g = \lambda_0/\sqrt{\epsilon_r}$  at the operating wavelength  $\lambda_0$ . On the thick substrates between  $0.037\lambda_g$ – $0.229\lambda_g$ , Chang *et al.* [1] and Kara [2] have obtained the resonance frequency of rectangular patches experimentally to achieve the bandwidth up to 22%. For such patches, the transmission line model and the cavity model compute the resonance frequency with average deviation between 3%–16% [1], [3], [4]. Usually, these experimental results have not been compared against any full-wave method. However, we have noted that the method of moments (MOM) based commercial software, Ensemble [5], provides deviation up to 12% for the patches of Chang *et al.* and up to 36% for the patches of Kara.

An improved form of the cavity model, called modified Wolff model (MWM), has been developed by Verma and co-workers [6], [7] to compute the resonance frequency of rectangular patches on the isotropic/anisotropic substrate. In this paper, apart from the effect of anisotropy in the substrate, we further account for the effect of all kind of losses in the patch on the resonance frequency. Thus, the present form of MWM computes accurately the resonance frequency of rectangular patches on electrically very thick substrate. The maximum deviation in the present model is 2% against the experimental results of Chang *et al.* and Kara.

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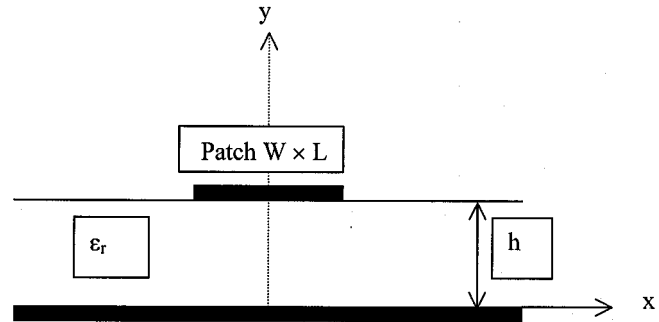


Fig. 1. Rectangular patch on thick substrate.

## II. IMPROVED CAVITY MODEL (MODIFIED WOLFF MODEL)

We replace the uniaxial anisotropic substrate by the equivalent isotropic dielectric layer of equivalent relative permittivity  $\epsilon_{req}$ , and equivalent substrate thickness  $h_{eq}$  [7]

$$\epsilon_{req} = \sqrt{\epsilon_{rx}\epsilon_{ry}}, \quad h_{eq} = h \sqrt{\frac{\epsilon_{rx}}{\epsilon_{ry}}} \quad (1)$$

where  $h$  is the actual thickness of substrate.

The above anisotropy relation has been obtained in the literature for the microstrip line on the anisotropic substrate by solving the boundary value problem [8]. The optic axis of the substrate is aligned to the geometrical axis of the patch.

The modified Wolff model (MWM) computes the resonance frequency of a rectangular patch on the lossy equivalent isotropic substrate [6], [7] by

$$f_r = \text{Re} \left[ \frac{v_0}{2\sqrt{\epsilon_{r dyn}^*}} \left[ \left( \frac{n}{L_{eff}} \right)^2 + \left( \frac{m}{W_{eff}} \right)^2 \right]^{\frac{1}{2}} \right] \quad (2)$$

In the normal cavity model,  $\epsilon_{r eff}^*$  or  $\epsilon_{r dyn}^*$  is taken as a real quantity. However, in the present case, the complex dynamic effective relative permittivity ( $\epsilon_{r dyn}^*$ ) takes care of the modal field variation given by the modal numbers  $m, n$ , and anisotropy in the substrate. It further accounts for the dielectric, conductor, surface-wave, and radiation losses. The  $\epsilon_{r dyn}^*$  is computed by [6], [7],

$$\epsilon_{r dyn}^* = \frac{C_{dyn}^* (\epsilon_{req}^*, h_{eq})}{C_{dyn}^* (\epsilon_{req}^* = 1, h_{eq})} \quad (3)$$

All above mentioned losses in the patch are accommodated by the concept of effective loss-tangent ( $\tan \delta_{eff}$ ). Thus, the complex equivalent relative permittivity of lossy substrate is  $\epsilon_{req}^*$  =

$\epsilon'_{req} - j\epsilon''_{req}$ , where  $\epsilon''_{req} = \epsilon'_{req} \tan \delta_{eff}$ . The  $\epsilon'_{req}$  for an anisotropic substrate is determined from (1).

The accurate computation of resonance frequency is done in three steps.

*Step I:* At first, we compute the approximate resonance frequency of a patch on the loss-less ( $\tan \delta_{eff} = 0$ ) anisotropic substrate by (1) and (2).

*Step II:* Next, we compute all  $Q$ -factors of the patch in order to determine  $\tan \delta_{eff}$ . The  $Q_d$ ,  $Q_c$ ,  $Q_{sw}$ , and  $Q_r$  i.e., the  $Q$ -factors, due to dielectric loss, conductor loss (in patch, ground plane and in the co-axial feed), surface-wave loss, and radiation loss, respectively, are determined by the following expressions:

$$Q_d = \frac{\pi \sqrt{\epsilon_{effav}}}{\lambda_0 \alpha_d}, \quad Q_c = \frac{\pi \sqrt{\epsilon_{effav}}}{\lambda_0 (\alpha_c + \alpha_{feed})},$$

$$Q_r = \frac{\pi}{4G_r Z (\epsilon_r, w/h)}, \quad Q_{sw} = \frac{Q_r \eta}{(1 - \eta)} \quad (4)$$

where

$$\epsilon_{effav} = \frac{\epsilon_{eff(w)} + \epsilon_{eff(L)}}{2}$$

and

$$\eta = \frac{P_{sp}}{P_{sp} + P_{sw}}. \quad (5)$$

Finally, we get  $\tan \delta_{eff}$  of the patch antenna

$$\tan \delta_{eff} = \frac{1}{Q_t} = \left[ \frac{1}{Q_d} + \frac{1}{Q_c} + \frac{1}{Q_{sw}} + \frac{1}{Q_r} \right]. \quad (6)$$

Once total  $Q$ -factor is known, we can also compute the 1:2 VSWR bandwidth of the patch antenna,

$$\%BW = \frac{100}{\sqrt{2}Q_t}. \quad (7)$$

The effective permittivities  $\epsilon_{eff(w)}$  and  $\epsilon_{eff(L)}$  have been determined from the width ( $w$ ) and length ( $L$ ) sides respectively by the variational method [6], [7]. The dielectric loss ( $\alpha_d$ ) is determined by using the standard expression and the conductor loss ( $\alpha_c$ ) in patch and ground plane is determined from the Wheeler's inductance rule [9]. The loss in the probe feed ( $\alpha_{feed}$ ) is determined by following Collin [10]. This loss is small and can be neglected. The radiation conductance ( $G_r$ ) of the radiating aperture is computed from the approximate expressions of James *et al.* [11]. The power in surface wave ( $P_{sw}$ ) and power in the radiated space-wave ( $P_{sp}$ ) are computed from the expressions of Pozar [12]. However, these expressions are valid for the infinite substrate which do not account for the standing wave inside the substrate.

*Step III:* The third step is to get the imaginary part of equivalent relative permittivity, i.e.,  $\epsilon''_{req}$  and then to compute the complex  $\epsilon_{req}^*$ . Equation (1) finally computes accurate resonance frequency of the patch. The bandwidth is also computed at the final resonance frequency.

### III. EXPERIMENTAL SETUP

Chang *et al.* and Kara have used different methods for measuring the resonance frequency of the rectangular patch on the thick PTFE substrates. Chang *et al.* have kept distance of the

TABLE I  
PATCHES OF CHANG *et al.* [1]: [NONWOVEN GLASS MICROFIBER REINFORCED SUBSTRATE, ANISOTROPY RATIO = 1.05,  $\epsilon_r = 2.33 \pm 0.02$ ,  $\tan \delta = 0.001$ ,  $t = 0.005$  mm]

Patch No.	W (mm)	L (mm)	h (mm)	h/ $\lambda_g$	Expt. [1] $f_r$ (GHz)	PCAAD [4] $f_r$ (GHz)	James [1] $f_r$ (GHz)	Ham. [1] $f_r$ (GHz)	MCM [12] $f_r$ (GHz)	Ensem. [5] $f_r$ (GHz)	MWM $f_r$ (GHz)
1	57.0	38.0	3.175	0.037	2.31	2.37	2.30	2.38	2.38	2.38	2.34
2	45.5	30.5	3.175	0.047	2.89	2.90	2.79	2.90	2.90	2.90	2.88
3	17.0	11.0	1.524	0.061	7.87	7.74	7.46	7.84	7.80	7.80	7.88
4	29.5	19.5	3.175	0.068	4.24	4.27	4.11	4.34	4.30	4.37	4.25
5	19.5	13.0	3.175	0.094	5.84	5.94	5.70	6.12	6.02	6.11	5.90
6	17.0	11.0	3.175	0.110	6.80	6.74	6.47	7.01	6.87	6.94	6.81
7	14.0	9.0	3.175	0.125	7.70	7.81	7.46	8.19	8.01	8.79	7.72
8	12.0	8.0	3.175	0.141	8.27	8.51	8.13	9.01	8.78	9.74	8.27
9	10.5	7.0	3.175	0.148	9.14	9.32	8.89	9.97	9.66	10.70	9.17
10	9.0	6.0	3.175	0.166	10.25	10.31	9.82	11.18	10.76	11.66	10.31
11	17.0	11.0	9.525	0.229	4.73	4.58	4.32	5.27	4.78	4.3	4.65
Maximum % Deviation :					3.17	8.67	11.42	6.17	17.78	1.69	
Average of absolute % Deviation :					1.63	3.62	5.35	2.96	7.81	0.54	
RMS % Deviation :					0.56	1.26	1.88	1.07	3.05	0.23	

TABLE II  
PATCHES OF KARA [2]: [WOVEN TYPE PTFE SUBSTRATE, ANISOTROPY RATIO = 1.2,  $\epsilon_r = 2.55 \pm 0.05$ ,  $\tan \delta = 0.002$ ,  $t = 0.005$  mm]

Patch No.	W (mm)	L (mm)	h (mm)	h/ $\lambda_g$	Expt. [2] $f_r$ (GHz)	PCAAD [4] $f_r$ (GHz)	MCM [14] $f_r$ (GHz)	Ensem. [5] $f_r$ (GHz)	MWM $f_r$ (GHz)
1	10.80	7.76	3.30	0.1405	8.00	8.409	9.010	9.697	7.835
2	14.50	9.87	4.50	0.1454	6.07	6.439	6.863	8.258	5.964
3	15.20	10.00	4.76	0.1475	5.82	6.249	6.638	7.940	5.778
4	12.55	7.90	4.00	0.1519	7.134	7.718	8.165	9.293	7.092
5	19.70	12.00	6.26	0.1553	4.66	5.015	5.215	5.030	4.606
6	14.40	8.14	4.76	0.1617	6.38	7.046	7.330	7.833	6.433
7	16.20	7.90	5.50	0.1754	5.99	6.685	6.980	5.950	6.060
8	27.56	12.56	9.52	0.1814	3.58	4.030	3.920	3.748	3.626
9	26.40	10.20	9.52	0.1976	3.90	4.439	3.938	3.895	3.956
10	26.20	9.74	9.52	0.2017	3.98	4.528	3.965	3.895	4.018
11	35.00	12.65	12.81	0.2032	2.98	3.416	2.958	2.920	3.031
12	23.00	7.83	8.45	0.2091	4.60	5.312	4.348	4.545	4.691
13	26.76	8.83	10.00	0.2119	3.98	4.580	3.800	3.814	4.027
14	34.00	10.80	12.81	0.2148	3.15	3.640	2.973	2.972	3.193
15	33.80	10.30	12.81	0.2182	3.20	3.705	2.980	3.005	3.242
16	31.30	9.20	12.00	0.2216	3.47	4.026	3.225	3.211	3.507
17	28.35	7.77	11.0	0.2284	3.90	4.519	3.388	3.563	3.909
Maximum % Deviation :					16.02	16.53	36.43	2.06	
Average of absolute % Deviation :					12.06	8.93	11.62	1.22	
rms % Deviation :					3.06	2.51	4.08	0.32	

probe (1.5 mm) fixed from edge of the patch. Thus, in their case, at the resonance frequency, the imaginary part of input impedance is not zero. They have obtained the resonance frequency of the patch corresponding to the maximum value of real part of the input impedance. Kara has measured the resonance frequency after experimentally optimizing the probe distance to get the return loss better than 30 dB. Both sets of experimental investigations do not mention about the anisotropy, loss tangent, and conductor thickness. We have taken this information from the handbook in [13]. Both the investigators have provided experimental results on the bandwidth also.

### IV. RESULTS AND CONCLUSION

Tables I and II show the experimental resonance frequency of the rectangular patches of Chang *et al.* [1] and Kara [2], respectively. The tables also contain the computed resonance fre-

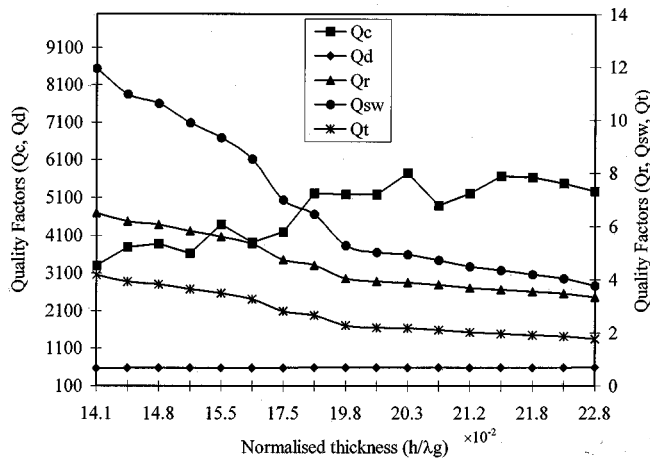


Fig. 2.  $Q$ -Factors due to conductor, dielectric, radiation, surface wave and total losses of patches of Kara shown in Table II.

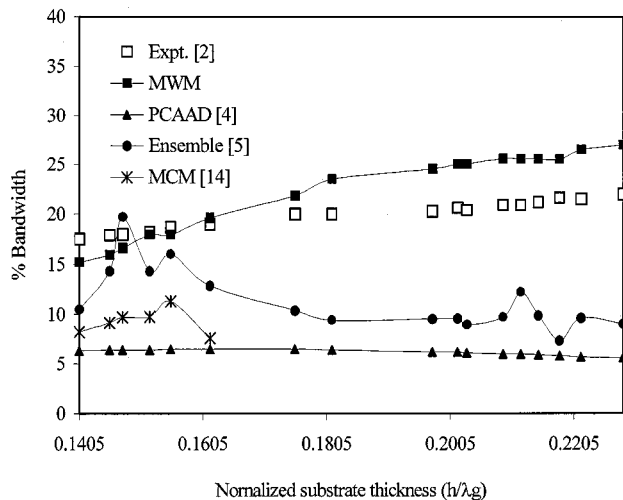


Fig. 3. % bandwidth with electrical thickness of patches in Table II.

quency by the cavity model of Richards *et al.*, as adopted by Pozar in his commercial software PCAAD [4], transmission line model of James [1], Hammerstad [1], multiport cavity model (MCM) of Benalla *et al.* [14], MOM of Ensemble [5], and, finally, the present version of the modified Wolff model (MWM).

It is obvious that the present form of the MWM is much better than the previous form of the cavity model, the transmission models, the multiport cavity model (MCM), and, finally, the MOM based commercial software, Ensemble. The MCM and Ensemble have deviation as high as 16% and 36%, respectively. The MWM has maximum deviation within 2% and average deviation 0.54%. Thus, the anisotropy and losses, mainly surface-wave and radiation losses, are key factors in influencing the resonance frequency of a patch radiator on the thick substrate which have not been accounted for in these models. It would be

interesting to compare these experimental results against the finite element method (FEM) and finite difference time domain (FDTD) based commercial software. However, we do not possess these software to do the same. In Figs. 2 and 3, we present the  $Q$ -factors and the bandwidth for the patches of Kara. Fig. 2 clearly shows the role of each  $Q$ -factor for a patch radiator on the thick substrate. The low  $Q$ -factor due to the surface wave loss decreases the efficiency to 60%, and it also degrades the radiation pattern. Finally, Fig. 3 compares the bandwidth computed by the above-mentioned models against the experimental bandwidth of Kara. The MWM shows better agreement with the experimental results as compared to models and software mentioned in this paper. For thickness above  $0.1605\lambda_g$ , the MCM does not work properly.

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